

Patent Application

COLLECTOR FOR EUV LIGHT SOURCE

Inventors

William N. Partlo
John Martin Algots
Gerry Blumenstock
Norbert Bowering
Alexander I. Ershov
Igor V. Fomenkov

And

Xiaojiang Pan

CYMER, Inc.
17075 Thornmint Court
San Diego, California 92127

"Express Mail" Label Number #EV 305 638 787 US

Date of Deposit March 10, 2004

I hereby certify that this paper or fee is being deposited with the United States Postal Service "Express Mail Post Office to Addressee" services under 37 C.F.R. 1.10 on the date indicated above and is addressed to the Assistant Commissioner for Patents, Washington, D.C. 20231.

Stephanie Sharrett

Typed Name of Person Mailing Paper or Fee


Signature

5
TITLE
Collector for EUV Light Source

10
William N. Partlo
J. Martin Algots
Gerry M. Blumenstock
Norbert R. Bowering
Alex I. Ershov
Igor V. Fomenkov
Xiaojiang Pan

15
FIELD OF THE INVENTION

The present invention relates to the field of the generation of EUV (soft-x-ray) light for such applications as semiconductor integrated circuit lithography exposure light sources, and more particularly to light collectors for such devices.

RELATED APPLICATIONS

This application is a continuation-in-part of U.S. Serial No. 10/409,254 filed April 8, 2003, the disclosure of which is incorporated by reference herein.

BACKGROUND OF THE INVENTION

25 The need for such applications as ever increasingly smaller critical dimensions for semiconductor integrated circuit manufacturing the need has arisen to move from the generation of Deep Ultraviolet ("DUV") light to Extreme Ultraviolet ("EUV") light, also referred to as soft-x-ray light. Various proposals exist for apparatus and methods for the generation of such light at effective energy levels to enable, e.g., adequate throughput in an EUV lithography tool (e.g., a
30 stepper scanner or scanner) over an acceptable lifetime between, e.g., replacements of major components.

Proposals exist for generating, e.g., light centered at a wavelength of 13.5nm using, e.g., Lithium which is introduced into and/or irradiated to form a plasma
35 which excites the lithium atoms to states from which decay results in large part in EUV light photons having an energy distribution centered about 13.5nm. The

plasma may be formed by an electrical discharge using a dense plasma focus electrode in the vicinity of a source of lithium in solid or liquid form, e.g., as discussed in United States Patent Nos. 6,586,757, entitled PLASMA FOCUS LIGHT SOURCE WITH ACTIVE BUFFER GAS CONTROL, issued to
5 Melynychuk et al. on July 1, 2003, and the above referenced patent application Ser. No. 10/409,254 filed April 8, 2003, and United States Patent No. 6,566,668, entitled PLASMA FOCUS LIGHT SOURCE WITH TANDEM ELLIPSOIDAL MIRROR UNITS, issued to Rauch et al. on May 20, 2003, and United States Patent No. 6,566,667, entitled PLASMA FOCUS LIGHT SOURCE WITH IMPROVED
10 PULSE POWER SYSTEM, issued to Partlo et al on May 20, 2003, which are assigned to the assignee of the present application and applications and patents and other references referenced therein, the disclosures of all of which are hereby incorporated by reference, and also other representative patents or published applications, e.g., United States Published Application No. 2002-0009176A1,
15 entitled X-RAY EXPOSURE APPARATUS, published on January 24, 2002, with inventors Amemlya et al. the disclosures of which are hereby incorporated by reference. In addition, as noted in, e.g., patents and published applications United States Patent No. 6,285,743, entitled METHOD AND APPARATUS FOR SOFT X-RAY GENERATOR, issued to Kondo et al. on September 4, 2001, United States
20 Patent No. 6,493,423, entitled METHOD OF GENERATING EXTREMELY SHORT-WAVE RADIATION ... , issued to Bisschops on December 10, 2002, United States Published Application 2002-0141536A1 entitled EUV, XUV AND X-RAY WAVELENGTH SOURCES CREATED FROM LASER PLASMA ..., Published on October 3, 2002, with inventor Richardson, United States Patent No.
25 6,377,651, entitled LASER PLASMA SOURCE FOR EXTREME ULTRAVIOLET LITHOGRAPHY USING WATER DROPLET TARGET, issued to Richardson et al. on April 23, 2002, United States Patent No. 6,307,913, entitled SHAPED SOURCE OF X-RAY, EXTREME ULTRAVIOLET AND ULTRAVIOLET RADIATION, issued to Foster et al. on October 23, 2001, the disclosures of which
30 are hereby incorporated by reference, the plasma may be induced by irradiating a target, e.g., a droplet of liquid metal, e.g., lithium or a droplet of other material

containing a target of, a metal, e.g., lithium within the droplet, in liquid or solid form, with, e.g., a laser focused on the target.

Since the amount of energy in the EUV light desired to be produced within the desired bandwidth, from the creation of such a plasma and resultant generation from the plasma of EUV light, is relatively enormous, e.g., 100 Watts/cm², its is
5 necessary to ensure that the efficiency of the collection of the EUV light be made as high as possible. It is also required that this efficiency not significantly deteriorate, i.e., be able to sustain such high efficiency, over relatively extended periods of operation, e.g., effectively a year of operation at very high pulse repetition rates
10 (4KHz and above) for an effective 100% duty cycle. Many challenges exist to being able to meet these goals aspects of which are dealt with in explaining aspects of the present invention regarding a collector for an EUV light source.

Some issues that are required to be addressed in a workable design include, e.g., Li diffusion into the layers of a multi layer normal angle of incidence reflecting
15 mirror, e.g., through an outer coating of ruthenium ("Ru"), with the multilayered mirror made, e.g., of alternating layers of Molybdenum ("Mo" or "Moly") and silicon ("Si") and the impact on, e.g., the primary and/or secondary collector lifetime; chemical reactions between, e.g., Li and Si and the impact on, e.g., the primary and/or secondary collector lifetimes; scatter of out of band radiation, e.g.,
20 from the laser producing the irradiation for ignition to form the plasma, e.g., 248 nm radiation from an KrF excimer laser required to be kept low to avoid any impact on resist exposure given that Deep UV resist types may be carried over into the EUV range of lithography and such out of band light scattered from the target can result in exposing the resist very efficiently; achieving a 100 W delivery of output light
25 energy to the intermediate focus; having a lifetime of a primary and secondary collector of at least 5G pulses; achieving the required conversion efficiency with a given source, e.g., a given target, e.g., a target droplet or target within a droplet, or other targets, the preservation of lifetime of the required multi layer mirrors at operational elevated temperatures and out of band radiation at center wavelengths
30 near, e.g., 13.5nm.

It is well known that that normal incidence of reflection ("NIR") mirrors can be constructed for wavelengths of interest in EUV, e.g., between about 5 and 20 nm, e.g., around 11.3 nm or 13.0-13.5 nm utilizing multi-layer reflection. The properties of such mirrors depend upon the composition, number, order, crystallinity, surface roughness, interdiffusion, period and thickness ratio, amount of annealing and the like for some or all of the layers involved and also, e.g., such things as whether or not diffusion barriers are used and what the material and thickness of the barrier layer is and its impact on the composition of the layers separated by the barrier layer, as discussed, e.g., in Braun, et al., "Multi-component EUV multi-layer mirrors, Proc. SPIE 5037 (2003) (Braun"); Feigl, et al., "Heat resistance of EUV multi-layer mirrors for long-time applications," *Microelectronic Engineering* 57-58, p. 3-8 (2001) ("Feigl"), United States Patent No. 6,396,900, entitled MULTILAYER FILMS WITH SHARP, STABLE INTERFACES FOR USE IN EUV AND SOFT X-RAY APPLICATION, issued to Barbee, Jr. et al. on May 28, 2002, based upon an application Ser. No. 10/847,744, filed on May 1, 2002 ("Barbee") and United States Patent No. 5,319,695, entitled MULTILAYER FILM REFLECTOR FOR SOFT X-RAYS, issued to Itoh et al. on June 7, 1994, based on an Application Ser. No. 45,763, filed on April 14, 1993, claiming priority to a Japanese application filed on April 21, 1992 ("Itoh").

Itoh discusses materials of different X-ray refractive indexes, for example, silicon (Si) and molybdenum (Mo), alternately deposited on a substrate to form a multilayer film composed of silicon and molybdenum layers and a hydrogenated interface layer formed between each pair of adjacent layers. Barbee discusses a thin layer of a third compound, e.g., boron carbide (B_4C), placed on both interfaces (Mo-on-Si and Si-on-Mo interface). This third layer comprises boron carbide and other carbon and boron based compounds characterized as having a low absorption in EUV wavelengths and soft X-ray wavelengths. Thus, a multi-layer film comprising alternating layers of Mo and Si includes a thin interlayer of boron carbide (e.g., B_4C) and/or boron based compounds between each layer. The interlayer changes the surface (interface) chemistry, which can result in an increase of the reflectance and increased thermal stability, e.g., for Mo/Si where inter-diffusion may be prevented

or reduced, resulting in these desired effects. Barbee also discusses varying the thickness of the third layer from the Mo-on-Si interface to the Si-on-Mo interface. Barbee also discusses the fact that typically the sharpness of the Mo-on-Si interface would be about 2.5 times worse than that of the Si-on-Mo interface; however, due to the deposition of the interlayer of B₄C in the Mo-on-Si interface, such interface sharpness is comparable to that of the Si-on-Mo interface. Braun discusses the use of carbon barrier layers to reduce inter-diffusion at the Mo-Si boundaries to improve the thermal stability and lower internal stress and at the same time increasing reflectivity. Braun notes that normally the Mo-Si boundary forms MoSi₂ at the interface in varying thicknesses at the Mo-on-Si boundary and the Si-on-Mo boundary, and also that the morphology of the Mo and/or Si layers can be influenced by barrier layers of, e.g., carbon content. In addition Braun notes the impact of barrier layer formation on interface roughness of the Mo-Si interface without a barrier layer. Braun reports a reflectance at $\lambda=13.3$ nm of 70.1% using Mo/SiC multi-layers. The reduction in internal stress using B₄C even with annealing as compared to Mo/Si/C multi-layers, which impacts the ability to use such multi-layer mirrors for curved mirrors is also discussed. Braun also discusses the tradeoff between interlayer contrast, impacting reflectivity, and absorption in the multi-layer configurations, such that, e.g., NbSi layers with lower absorption in the Nb but also lower contrast, and Ru/Si with higher contrast but also higher absorption in the Ru layer, both performing less effectively than a Mo/Si multi-layer stack. Braun also discusses the theoretical utility of using three layers of, e.g., Mo/Si/Ag or Mo/Si/Ru, which have theoretically higher reflectivity, but that the Ag embodiment fails to achieve the theoretical reflectivity due to voids in the Ag layer at desired thicknesses and a calculated best reflectivity of a Mo/Si/C/Ru multi-layer stack at $\lambda=13.5$ nm, with a thickness constrained in the Mo layer to prevent crystallization in the Mo layer. However, Braun also finds that the Mo/Si/C/Ru multi-layer stacks do not live up to theoretical calculated reflectivity expectation, probably due to an initial Mo layer deposition surface roughness that propagates upward through the stack. Feigl discusses the impact of elevated temperatures up to 500°C on the structural stability of, e.g., Mo/Si and Mo/Mo₂C/Si/Mo₂C multilayer stacks, including the use of

ultrathin Mo₂C barrier layers. Feigl notes that the barrier layer prevents the formation of inter-diffusion layers of MoSi_x due to annealing of the Mo and Si at temperatures above, e.g., 200°C and that Mo/Mo₂C/Si/Mo₂C and Mo₂C/Si systems remain stable up to 600°C. The former system having ultrathin Mo₂C barrier layers
5 (MoSi₂ is also suggested but not tested) layers and the latter is formed by substituting Mo₂C for Mo in a multilayer system. The reflectivity of the Mo₂C/Si system remained above .8 through 600°C according to Feigl, whereas the Mo/Mo₂C/Si/Mo₂C system tailed off to slightly less than .7 at that temperature, and even decreased to about .7 at 400°C.

10 Applicants in the present application propose certain other materials for barrier layers and other potential improvements to the multi-layer stack for EUV applications.

SUMMARY OF THE INVENTION

15 A method and apparatus for debris removal from a reflecting surface of an EUV collector in an EUV light source is disclosed which may comprise the reflecting surface comprises a first material and the debris comprises a second material and/or compounds of the second material, the system and method may comprise a controlled sputtering ion source which may comprise a gas comprising
20 the atoms of the sputtering ion material; and a stimulating mechanism exciting the atoms of the sputtering ion material into an ionized state, the ionized state being selected to have a distribution around a selected energy peak that has a high probability of sputtering the second material and a very low probability of sputtering the first material. The stimulating mechanism may comprise an RF or microwave
25 induction mechanism. The gas is maintained at a pressure that in part determines the selected energy peak and the stimulating mechanism may create an influx of ions of the sputtering ion material that creates a sputter density of atoms of the second material from the reflector surface that equals or exceeds the influx rate of the plasma debris atoms of the second material. A sputtering rate may be selected for a
30 given desired life of the reflecting surface. The reflecting surface may be capped. The collector may comprise an elliptical mirror and a debris shield which may

comprise radially extending channels. The first material may be molybdenum, the second lithium and the ion material may be helium. The system may have a heater to evaporate the second material from the reflecting surface. The stimulating mechanism may be connected to the reflecting surface between ignition times. The reflecting surface may have barrier layers. The collector may be a spherical mirror in combination with grazing angle of incidence reflector shells, which may act as a spectral filter by selection of the layer material for multi-layer stacks on the reflector shells. The sputtering can be in combination with heating, the latter removing the lithium and the former removing compounds of lithium, and the sputtering may be by ions produced in the plasma rather than excited gas atoms.

BRIEF DESCRIPTION OF THE DRAWINGS

FIG. 1 shows a schematic view of an overall broad conception for a laser-produced plasma EUV light source according to an aspect of the present invention;

FIG. 1A shows schematically the operation of the system controller according to an aspect of an embodiment of the present invention;

FIG. 2A shows a side view of an embodiment of an EUV light collector according to an aspect of the present invention looking from an irradiation ignition point toward an embodiment of a collector according to an embodiment of the present invention;

FIG. 2B shows a cross-sectional view of the embodiment of FIG. 2A along the lines 2B in FIG. 2A;

FIG. 3 shows an alternative embodiment of a normal angle of incidence collector according to an aspect of the present invention;

FIG. 4 shows a schematic view of a normal angle of incidence collector debris management system according to an aspect of the present invention;

FIG.'s 5a-c show a timing of the provision of a collector cleaning signal/current at RF and/or DC to the collector mirror according to an aspect of an embodiment of the present invention;

FIG.'s 6a and b show schematic views in cross section of aspects of embodiments of the present invention relating to grazing angle of incidence collectors;

FIG. 7 shows a plot of grazing angle of incidence reflectivity for a variety of reflective surfaces at given wavelengths of relevance at an angle of incidence of 5 degrees;

FIG. 8 shows a plot of grazing angles of incidence reflectivity for a variety of reflective surfaces at given wavelengths of relevance for 15 degrees;

FIG. 9 shows a schematic view of an alternative embodiment of a collector according to an aspect of the present invention;

FIG. 10 shows a calculated number of lithium atoms per droplet vs. droplet diameter, useful in illustrating an aspect of an embodiment of the present invention;

FIG. 11 shows a calculated influx of lithium atoms onto a mirror surface vs. mirror radius useful in illustrating an aspect of an embodiment of the present invention;

FIG. 12 shows a calculated required lithium thickness sputter rate vs. mirror diameter useful in illustrating an aspect of an embodiment of the present invention.

FIG. 13 shows a required ratio of molybdenum sputter rate to lithium sputter rate vs. mirror radius in order to have a 1-year life with a 300 pair multi-layer coated mirror useful in illustrating an aspect of an embodiment of the present invention;

FIG. 14 shows sputter yield for lithium, silicon, and molybdenum with helium ions useful in illustrating an aspect of an embodiment of the present invention;

FIG. 15 shows normalized helium ion energy along with sputter yields for lithium, silicon, and molybdenum useful in illustrating an aspect of an embodiment of the present invention;

FIG. 16 shows helium ion current density along with the sputter yield of lithium, silicon, and molybdenum useful in illustrating an aspect of an embodiment of the present invention;

FIG. 17 shows total helium ion sputter rates for lithium, silicon, and molybdenum useful in illustrating an aspect of an embodiment of the present invention;

FIG. 18 shows normalized lithium ion energy along with sputter yields for lithium and molybdenum useful in illustrating an aspect of an embodiment of the present invention;

FIG. 19 shows a radiated power density vs. temperature for a black body useful in illustrating an aspect of an embodiment of the present invention;

FIG. 20 shows a schematic view of an aspect of an embodiment of the present invention;

FIG.'s 21 A and B show results of experiments regarding the stopping power of helium and argon buffer gases against both tin and lithium ions according to aspects of an embodiment of the present invention; and,

FIG.'s 22A-E show results of further examination of the stopping power of helium and argon buffer gases against both lithium and tin according to aspects of an embodiment of the present invention.

DETAILED DESCRIPTION OF PREFERRED EMBODIMENTS

Turning now to FIG. 1 there is shown a schematic view of an overall broad conception for an EUV light source, e.g., a laser produced plasma EUV light source 20 according to an aspect of the present invention. The light source 20 may contain a pulsed laser system 22, e.g., a gas discharge laser, e.g., an excimer gas discharge laser, e.g., a KrF or ArF laser operating at high power and high pulse repetition rate and may be a MOPA configured laser system, e.g., as shown in United States Patents Nos. 6,625,191, 6,549,551, and 6,567,450. The laser may also be, e.g., a solid state laser, e.g., a YAG laser. The light source 20 may also include a target delivery system 24, e.g., delivering targets in the form of liquid droplets, solid particles or solid particles contained within liquid droplets. The targets may be delivered by the target delivery system 24, e.g., into the interior of a chamber 26 to an irradiation site 28, otherwise known as an ignition site or the sight of the fire ball. Embodiments of the target delivery system 24 are described in more detail below.

Laser pulses delivered from the pulsed laser system 22 along a laser optical axis 55 through a window (not shown) in the chamber 26 to the irradiation site, suitably focused, as discussed in more detail below in coordination with the arrival of a target produced by the target delivery system 24 to create an ignition or fire ball that forms an x-ray (or soft x-ray (EUV) releasing plasma, having certain characteristics, including wavelength of the x-ray light produced, type and amount of debris released from the plasma during or after ignition, according to the material of the target.

The light source may also include a collector 30. e.g., a reflector, e.g., in the form of a truncated ellipse, with an aperture for the laser light to enter to the ignition site 28. Embodiments of the collector system are described in more detail below. The collector 30 may be, e.g., an elliptical mirror that has a first focus at the ignition site 28 and a second focus at the so-called intermediate point 40 (also called the intermediate focus 40) where the EUV light is output from the light source and input to, e.g., an integrated circuit lithography tool (not shown). The system 20 may also include a target position detection system 42. The pulsed system 22 may include, e.g., a master oscillator-power amplifier ("MOPA") configured dual chambered gas discharge laser system having, e.g., an oscillator laser system 44 and an amplifier laser system 48, with, e.g., a magnetic reactor-switched pulse compression and timing circuit 50 for the oscillator laser system 44 and a magnetic reactor-switched pulse compression and timing circuit 52 for the amplifier laser system 48, along with a pulse power timing monitoring system 54 for the oscillator laser system 44 and a pulse power timing monitoring system 56 for the amplifier laser system 48. The pulse power system may include power for creating laser output from, e.g., a YAG laser. The system 20 may also include an EUV light source controller system 60, which may also include, e.g., a target position detection feedback system 62 and a firing control system 65, along with, e.g., a laser beam positioning system 66.

The target position detection system may include a plurality of droplet imagers 70, 72 and 74 that provide input relative to the position of a target droplet, e.g., relative to the ignition site and provide these inputs to the target position detection feedback system, which can, e.g., compute a target position and trajectory, from which a target error can be computed, if not on a droplet by droplet basis then on average, which is then provided as an input to the system controller 60, which can, e.g., provide a laser position and direction correction signal, e.g., to the laser beam positioning system 66 that the laser beam positioning system can use, e.g., to control the position and direction of the laser position and direction changer 68, e.g., to change the focus point of the laser beam to a different ignition point 28.

The imager 72 may, e.g., be aimed along an imaging line 75, e.g., aligned with a desired trajectory path of a target droplet 94 from the target delivery

mechanism 92 to the desired ignition site 28 and the imagers 74 and 76 may, e.g., be aimed along intersecting imaging lines 76 and 78 that intersect, e.g., along the desired trajectory path at some point 80 along the path before the desired ignition site 28.

5 The target delivery control system 90, in response to a signal from the system controller 60 may, e.g., modify the release point of the target droplets 94 as released by the target delivery mechanism 92 to correct for errors in the target droplets arriving at the desired ignition site 28.

10 An EUV light source detector 100 at or near the intermediate focus 40 may also provide feedback to the system controller 60 that can be, e.g., indicative of the errors in such things as the timing and focus of the laser pulses to properly intercept the target droplets in the right place and time for effective and efficient LPP EUV light production.

15 Turning now to FIG. 1A there is shown schematically further details of a controller system 60 and the associated monitoring and control systems, 62, 64 and 66 as shown in FIG. 1. The controller may receive, e.g., a plurality of position signal 134, 136 a trajectory signal 136 from the target position detection feedback system, e.g., correlated to a system clock signal provided by a system clock 116 to the system components over a clock bus 115. The controller 60 may have a pre-arrival tracking and timing system 110 which can, e.g., compute the actual position of the target at some point in system time and a target trajectory computation system 112, which can, e.g., compute the actual trajectory of a target drop at some system time, and an irradiation site temporal and spatial error computation system 114, that can, e.g., compute a temporal and a spatial error signal compared to some desired point in space and time for ignition to occur.

25 The controller 60 may then, e.g., provide the temporal error signal 140 to the firing control system 64 and the spatial error signal 138 to the laser beam positioning system 66. The firing control system may compute and provide to a resonance charger portion 118 of the oscillator laser 44 magnetic reactor-switched pulse compression and timing circuit 50 a resonant charger initiation signal 122 and may provide, e.g., to a resonance charger portion 120 of the PA magnetic reactor-

switched pulse compression and timing circuit 52 a resonant charger initiation signal, which may both be the same signal, and may provide to a compression circuit portion 126 of the oscillator laser 44 magnetic reactor-switched pulse compression and timing circuit 50 a trigger signal 130 and to a compression circuit portion 128 of the amplifier laser system 48 magnetic reactor-switched pulse compression and timing circuit 52 a trigger signal 132, which may not be the same signal and may be computed in part from the temporal error signal 140 and from inputs from the light out detection apparatus 54 and 56, respectively for the oscillator laser system and the amplifier laser system.

The spatial error signal may be provided to the laser beam position and direction control system 66, which may provide, e.g., a firing point signal and a line of sight signal to the laser beam positioner which may, e.g. position the laser to change the focus point for the ignition site 28 by changing either or both of the position of the output of the laser system amplifier laser 48 at time of fire and the aiming direction of the laser output beam.

Turning now to FIG.'s 2A and 2B there is shown, respectively a schematic view side view of a collector 30 looking into the collector mirror 150, and a cross-sectional view of the rotationally symmetric collector mirror 150 arrangement along cross-sectional lines 2B in FIG. 2A (although the cross-sectional view would be the same along any radial axis in FIG. 2A).

As shown in FIG. 2A the elliptical collection mirror 150 is circular in cross section looking at the mirror, which may be the cross-section at the greatest extension of the mirror, which is shown in FIG. 1A to be almost to the focus point 28 of the elliptical mirror 150, so as not to block target droplets 94 from reaching the ignition point designed to be at the focus point 28. It will be understood, however, that the mirror may extend further towards the intermediate focus, with a suitable hole in the mirror (not shown) to allow passage of the target droplets to the focus point. The elliptical mirror may also have an aperture 152, e.g., shown to be circular in FIG. 2A, to allow entry of the LPP laser beam 154, e.g., focused through focusing optics 156, through the mirror 150 to the ignition point 28 desired to be at the focus of the elliptical mirror. The aperture 152 can also be, e.g., more tailored to the beam

profile, e.g., generally rectangular, within the requirements, if any of modifying the beam optical path to make corrections of the focus of the laser beam 154 on an ignition site, depending upon the type of control system employed.

Also shown in FIG.'s 2A and 2B is a debris shield 180 according to an aspect of an embodiment of the present invention. The debris shield 180 may be made up of a plurality of thin plates 182, made, e.g., of thin foils of molybdenum, extending radially outward from the desired ignition site and defining narrow planar radially extending channels 184 through the debris shield 180. The illustration of FIG. 2A is very schematic and not to scale and in reality the channels are as thin as can possibly be made. Preferably the foil plates 182 can be made to be even thinner than the channels 184, to block as little of the x-ray light emitted from the plasma formed by ignition of a target droplet 94 by the laser beam 155 focused on the ignition site 28.

Seen in cross section in FIG. 2B, the functioning of the channels 182 in the debris shield 180 can be seen. A single radial channel is seen in FIG. 2B and the same would be seen in any section of the collector 30 through the rotationally symmetric axis of rotation of the collector mirror 150 and debris shield 180 within a channel of the debris shield 180. Each ray 190 of EUV light (and other light energy) emitted from the ignition site 28 traveling radially outward from the ignition site 28 will pass through a respective channel 182 in the debris shield 180, which as shown in FIG. 2B may, if desired, extend all the way to the collection mirror 150 reflective surface. Upon striking the surface of the elliptical mirror 150, at any angle of incidence, the ray 190 will be reflected back within the same channel 180 as a reflected ray 192 focused on the intermediate focus 40 shown in FIG. 1.

Turning now to FIG. 3 there is shown an alternative embodiment according to an aspect of an embodiment of the present invention. In this embodiment, the debris shield 180 is not shown for simplicity and this embodiment can be utilized with or without a debris shield as appropriate, as discussed in more detail below, as can also, e.g., the single elliptical collector mirror shown in FIG.'s 2A and B. In this embodiment a secondary collector reflecting mirror 200 has been added, which may comprise, e.g., a section of a spherical mirror 202, having a center at the ignition site

28, i.e., the focus of the elliptical mirror 150, and with an aperture 210 for the passage of the light from the collector mirror 150 to the intermediate focus 40 (shown in FIG. 1). The collector mirror 150 functions as discussed above in regard to FIG.'s 2A and 2B with respect to rays 190 emitted from the ignition point 28 toward the collector mirror 150. Rays of light 204 emitted from the ignition site 28 away from the collector mirror 150 which strike the section of the spherical mirror 202, will be reflected back through the focus of the elliptical collector mirror 132 and, pass on to the elliptical collector mirror 150 as if emitted from the focus 28 of the elliptical mirror 150, and, therefore also be focused to the intermediate focus 40. It will be apparent that this will occur with or without the presence of the debris shield 180 as described in relation to FIG.'s 2A and 2B.

Turning now to FIG. 4 there is shown schematically another aspect of debris management according to an embodiment of the present invention. FIG. 4 shows a collector mirror 150 connected to a source of current, e.g., DC voltage source 220. This current can be, e.g., one embodiment of the present invention in which the current maintains the reflector at a selected temperature to, e.g., evaporate deposited lithium. An alternate concept for lithium removal from the first collector mirror is to employ helium ion or hydrogen ion sputtering. The low mass of these ions, when kept at low energies ($<50\text{eV}$), e.g., can lead to extremely low sputter yield for, e.g., the molybdenum layer and/or the silicon layer, e.g., in an EUV multi-layer mirror fabricated with Mo/Si layers.

Turning now to FIG. 4 there is shown a debris cleaning arrangement according to an aspect of an embodiment of the present invention. As shown in FIG. 4, a source of current, e.g., DC voltage source 220 may, e.g., be connected to the collector mirror 150, e.g., to a metal, e.g., aluminum or nickel backing (not shown) for the mirror 150. The mirror 150 may thus be heated to an elevated temperature above that of the surrounding gas, e.g., helium gas, making up the content of the EUV light source chamber 26 interior. Other heating of the reflector may occur according to alternative embodiments of the invention, e.g., by radiant heating from, e.g., a heat lamp (not shown) in the vessel 26.

Another aspect of debris cleaning may incorporate, e.g., as shown in FIG. 4, e.g., the introduction of RF, e.g., from a source of RF frequency voltage 230 and an antenna, shown schematically at 232 within the chamber 26 in FIG. 4. In fact, the RF, as with the DC shown in FIG. 4, may be connected to the mirror 150 or a metallic backing (not shown) in which event a dark shield (not shown) made of a suitable conductive material and connected to ground potential may be formed over the back of the collector mirror 150, separated from the mirror 150 by an insulator, e.g., an air gap, and the potential, e.g., DC from DC source 220 that is connected also to the mirror 150.

As shown in FIG.'s 5a-c, for a given periodic LPP ignition at times t_1 , t_2 , t_3 , the RF may be replaced by a DC voltage during the time when the ignitions occur at t_1 , t_2 , t_3 , and for a short time on either side of the ignition time, with RF between such times, at least directly after the ignition, if not completely through the next occurrence of the DC potential during the next ignition. Also shown is that the DC from source 220 may be a positive potential during the time of the respective ignition, perhaps coextensive with the continuous voltage from the RF source 230, and a negative potential between such positive pulses.

The voltage applied to the collector mirror 150 is meant to, on the one hand, evaporate metallic debris, e.g., lithium emitted from the plasma during and after ignition of a target droplet of such lithium or other target metallic material. Also evaporated could be metallic elements such as K, Fe, Na or the like that appear due, e.g., to impurities in the lithium target droplets themselves and are similarly deposited on the collector mirror 150 surface after ignition.

The RF is meant to form a localized ionic plasma, e.g., of excited He atoms in the vicinity of the collector mirror 150 surface, with the intent that these excited ions in the localized plasma may strike lithium atoms or compounds of lithium on the collector mirror 150 and sputter them off of the mirror surface. This embodiment of the invention contemplates, e.g., a balancing between the evaporation mechanism and the sputtering mechanism, e.g., if the RF is at <500W power (at 13.65MHz, as dictated by federal regulations for RF frequency sputtering) then the mirror temperature should be maintained at or near some desired

temperature and if the RF is increased, e.g., to >500 W at 13.65 MHz then the temperature can correspondingly be reduced.

Turning now to FIG.'s 6A and B there are shown aspects of embodiments of the present invention relating to alternative collector arrangements. As shown in FIG.'s 6A and B a collector 225 may be composed of, e.g., a plurality of nested shells, forming, e.g., different sections made up, e.g., of elliptical and parabolic reflecting shells, e.g., in FIG. 6a parabolic shells 230 and 240 and elliptical shells 250 and 260. The elliptical shells, e.g., 230 and 240 may be comprised of respectively first parabolic reflecting surfaces 233, 242, and second parabolic reflecting surfaces 234, 244. The elliptical sections 250 and 260 may be comprised of, e.g., elliptical reflecting surfaces 252 and 262. In FIG. 6B there is shown an alternative embodiment with an additional two parabolic shell sections 232 and 236, with section 232 comprising, e.g., a first parabolic reflecting surface 231 and a second parabolic reflecting surface 234, and section 236 comprising, e.g., a first parabolic reflecting surface 237, a second parabolic reflecting surface 238 and a third parabolic reflecting surface 239.

Each of the reflecting shells 230, 240, 250 and 260 are arranged to reflect between them 100 percent of the light emitted from the ignition point 21 within a section of a sphere from 11° to 55° from an axis of rotation 310 generally aligned with the focus of the collector 225 reflecting shells, with the shells 230, 240, 250 and 260 being generally symmetric about this axis of rotation 310 also. By way of example, the embodiment of FIG. 6A shows an embodiment where essentially all of the light in the portion of the sphere just described enters at least one of the shells 230, 240, 250 and 260. In the case of parabolic shell sections 230 and 240 is incident on the first reflecting surfaces 233, 242, and either reflected towards the intermediate focus 40, or is also reflected off of the respective second reflective surface 234, 244 to the intermediate focus. In the case of the elliptical shell sections 250, 260 all of the light entering each such shell 250, 260 is reflected to the intermediate focus, e.g., because the ellipses formed by the reflecting surfaces 252, 262 each have a first focus at the ignition point 28 and a second focus at the intermediate focus 40.

Depending on the material of the respective reflecting surfaces 233, 234, 242, 244, 252 and 262, the angle of incidence of the particular rays, the number of reflections in a given shell section 230, 240, 250 and 260, a certain average efficiency of reflection will occur and also depending on the construction of the shells a certain percentage of the available light will enter each section 230, 240, 250 and 260, such that, as illustrated in FIG. 6A 19% is reflected and focused in shell section 230 at an average total efficiency of 65%, 17% is reflected and focused in shell section 240 at an average total efficiency of 75%, 43% is reflected in shell section 250 at an average total efficiency of 80% and 21% is reflected and focused in shell section 260 with an average total efficiency of 91%.

FIG. 6B shows an alternative embodiment adding two more parabolic shell sections 232, 236. These added sections may serve, e.g., to collect more light up to about 85% from the axis of rotation and at least one of the added sections may have a first reflective surface 237, a second reflective surface 238 and a third reflective surface 239. As can be seen from FIG. 6B, e.g., a ray 290 of emitted light from the source or ignition point may just enter the parabolic reflective shell section 236 and be reflected as ray 292 to the second reflective surface 238 and then reflected as ray 294 to the third reflective surface 239 and then form focused ray 296. Similarly a ray 300 may just enter the parabolic reflector shell 236 at the other extremity of the shall opening and also be reflected off of the first reflective surface 237 as ray 320 and the second reflective surface 237 as ray 304 and the very end of the third reflective surface as focused ray 306. In the case of, e.g., one of the parabolic shell sections, e.g., section 240 a ray 280 may just enter this section 240 and be reflected off of the first parabolic reflecting surface 242 as ray 282 and the very end of the second parabolic reflecting surface as focused ray 283, and another ray 284 may just enter the section 240 to be reflected off of the second reflective surface 244 as focused ray 286. In the case of one of the elliptical shell sections, e.g., 250, a ray 308 emitted from the ignition point may just enter the shell section 250 and be reflected off of the elliptical reflecting surface 252 as a focused ray 309, and a ray 318 just enter the shell section 250 at the opposite side as ray 308 and be reflected as focused ray 319.

Turning now to FIG.'s 7 and 8 there is shown a plot of grazing angle of incidence reflectivity for (1) a single layer Ruthenium reflecting surface and (2) a Mo/Si bilayer stack with a Mo/Si 14nm thick single Mo layer and a 4nm single Si layer and (3) a ten period multi-layer Mo/Si stack, with a pitch of 9.4nm and a MO/Si thickness ratio of 22.5:1, having, e.g., 40 multi-layer stacks, each for grazing angles of incidence of 5° and 15°. In each of the reflectors with Mo/Si a molybdenum substrate is assumed. In the case where spectral purity is a part of the specification for the delivered light, collectors can be tuned to a certain wavelength, with some given bandwidth spread, e.g., by using the reflective properties of, e.g., a nested shell collector to favor reflectivity near a selected center wavelength, e.g., in the embodiments of FIG. 6A and B.

FIG. 9 shows aspects of an embodiment of the present invention. In this embodiment a collector assembly 330 may comprise, e.g., a portion of a spherical mirror reflecting surface 332, which may be a normal angle of incidence multi-layer stack, reflecting the light produced from the ignition point 28 to one of, e.g., three nested elliptical shell sections 336, 338 and 340 in a nested elliptical shell collector 334. Each of the shell sections 336, 338 and 340 may have a reflecting surface 366, 368, 369 that is on the inside of a respective shell 360, 362, 364. As shown in FIG. 9, e.g., the shell section 336 may receive the light from a rim section 370 of the spherical mirror 332, the shell section 338 may receive the light from an intermediate section of the spherical mirror 332 and the shell section 340 receives the light reflected from a central portion of the spherical mirror 332.

The shell sections 336, 338 and 340 may be coated with a multi-layer of Mo/Si rather than the conventionally proposed thick single layer of Ru. According to aspects of an embodiment of the present invention two reflections occur, e.g., one from the spherical mirror and one in each shell, e.g., for shells having elliptical reflecting surfaces, at grazing angles between about 5° and 15°, as can be seen from FIG's 7 and 8. This can, e.g., significantly reduce, e.g., a significant amount of out of band EUV radiation, e.g., assuming that 13.5 is the desired band. Ru mirrors, e.g., in a Wolter-type configuration remain very reflective for both 13.5nm and 11nm at both 5° and 15° grazing angles of incidence, whereas Mo/Si stacks of

grazing angle of incidence reflective coatings, as shown in FIG.'s 7 and 8 can be much more selective, especially around 15°.

The above described embodiment does not have the spatial purity of, e.g., a grating spectral purity filter, as has been proposed in the art, but it does have a significant advantage in reflectivity and preservation of in-band EUV radiation over the other solutions, e.g., a grating filter, proposed in the art.

A lithium LPP EUV light source according to aspects of embodiments of the present invention, could employ a solid stream of liquid lithium or a lithium droplet source. For a droplet source, the number of atoms per droplet can be calculated and for a solid stream one can assume that only material within the focused beam constitutes a droplet at ignition, although, from a debris standpoint adjacent material in the stream may also form debris, particularly if struck by lower energy laser radiation in the skirts of the energy distribution of the focused laser beam.

Since it is contemplated that it is desirable for the droplet source to have a droplet size matched to the focused beam, both types of target source can be considered to have the same droplet size given by a droplet diameter, d_{droplet} . The volume of the droplet is then given by:

$$V_{\text{droplet}} = \frac{1}{6} \pi d_{\text{droplet}}^3 \{1\}.$$

Calculating the number of atoms per droplet follows from the density of, e.g., lithium and its atomic weight. The mass of the droplet is:

$$M_{\text{droplet}} = V_{\text{droplet}} \rho_{\text{lithium}} \{2\};$$

where $\rho_{\text{lithium}} = 0.535 \text{ g/cm}^3$ is the density of lithium, such that:

$$M_{\text{droplet}} = 0.280 \cdot d_{\text{droplet}}^3 \{3\};$$

where the droplet diameter is in centimeters and the resulting mass is in grams. The number of atoms in the droplet is then given by dividing the droplet mass by the atomic mass of lithium and converting units properly:

$$N_{\text{atoms}} = \frac{M_{\text{droplet}} (\text{g})}{M_{\text{lithium atom}} (\text{amu})} \cdot \frac{1 \text{ amu}}{1.6605 \times 10^{-24} \text{ g}} \{4\};$$

where $M_{\text{lithium atom}} = 6.941 \text{ amu}$, i.e.,

$$N_{\text{atoms}} = 2.43 \times 10^{22} \cdot d_{\text{droplet}}^3 \{5\};$$

where the diameter of the droplet is in centimeters. Converting the droplet diameter from centimeters to micrometers gives:

$$N_{\text{atoms}} = 2.43 \times 10^{10} \cdot d_{\text{droplet}}^3 \{6\}.$$

The number of atoms per droplet versus droplet size is shown in FIG. 10
 5 Also shown in FIG. 10 is the number of 13.5nm photons contained in, e.g., a single 40mj pulse. The 40mj pulse example assumes a 10% conversion efficiency into 4π steradians and a 400mj laser pulse. The number of 13.5nm photons per pulse is given by:

$$10 \quad N_{\text{Photons}} = \frac{13.5\text{nm OpticalPulseEnergy(mJ)}}{E_{\text{Photon}}(\text{eV}) \cdot 1.6 \times 10^{-16}(\text{mJ/eV})} \{7\};$$

where the 13.5nm photon energy is 91.6eV. The resulting number of photons for a 40mj pulse is 2.72×10^{15} . A, e.g., 50 μm droplet has one lithium atom for every 13.5nm photon. Normally one could assume multiple photons emitted from each
 15 emission element. This assumption would allow use of a smaller droplet diameter than 50 μm . A smaller droplet diameter can be important because the lithium usage and lithium deposition rates, e.g., on the collector optics, scale as the cube of the droplet diameter.

Assuming that there is no lithium recovery, according to a possible aspect of
 20 an embodiment of the present invention, then calculating, e.g., the yearly usage of lithium is given by the number of pulses per year times the amount per pulse. Assuming, by way of example a repetition rate, RR, and a duty cycle, DC, the resulting mass usage is, e.g.,:

$$25 \quad \text{Mass Per Year} = M_{\text{droplet}} \cdot \text{RR} \cdot 60\text{sec/min} \cdot 60\text{min/hr} \cdot 24\text{hr/day} \cdot 365\text{day/yr} \cdot \text{DC} \{8\}.$$

i.e.,

$$\text{Mass Per Year} = 8.83 \times 10^{-6} \cdot d_{\text{droplet}}^3 \cdot \text{RR} \cdot \text{DC} \{9\};$$

where the droplet diameter is in micrometers and the resulting mass is in grams. For example, a system with no lithium recovery running at 6kHz with a droplet diameter

of 50 μ m running at 100% duty cycle for a full year would consume 6,622 grams or about a 12.3 liter volume of lithium. A droplet diameter of 25 μ m under similar conditions would consume only 828 grams or about 1.5 liters of lithium.

Assuming that the lithium droplet, once heated by the laser pulse, expands in all directions uniformly, the atomic flux will fall off as the square of the distance from the laser-droplet interaction point (ignition site). The number of atoms emitted from the interaction point per second is the number of atoms per droplet times the repetition rate:

$$\text{Total Atomic Emission} = 2.43 \times 10^{10} \cdot d_{\text{droplet}}^3 \cdot \text{RR} \quad \{10\};$$

where the droplet diameter is in micrometers and RR is the laser repetition rate in Hz.

The atomic flux (atoms/cm²) through the surface an imaginary sphere centered at the ignition site will be the total atomic emission divided by the surface area in centimeters:

$$\text{Atomic Flux} = 1.93 \times 10^9 \frac{d_{\text{droplet}}^3 \cdot \text{RR}}{r_{\text{sphere}}^2} \quad \{11\};$$

The resulting flux is in units of atoms/cm²s. FIG. 11 shows the rate of lithium influx onto the mirror surface vs. mirror radius for several droplet diameters, i.e., (1) 25 μ m. (2) 55 μ m, (3) 100 μ m and (4) 200 μ m, assuming a 6kHz repetition rate and 100% duty cycle.

In order to maintain high mirror reflectivity, the influx of lithium onto the mirror surface can, e.g., be exceeded by the sputter rate of lithium, e.g., caused by incident helium ions. In addition, for long mirror lifetime the sputter rate of molybdenum by these same, e.g., helium ions must then be many orders of magnitude slower than that for, e.g., lithium.

The required ratio of sputter rate of the first and second metals, e.g., molybdenum to lithium, in order to achieve, e.g., a 1 year lifetime for the multi-layer coated collector mirror can be calculated, e.g., by assuming use of, e.g., a multi-layer stack with 300 layer pairs, e.g., so that erosion of, e.g., the first 200 layer pairs leaves a still comfortably effective 100 good pairs, i.e., still maintaining high

reflectivity. Also assumed is a sputter rate for the silicon layers that is much higher than that for the first metal, e.g., molybdenum layers and thus provides a negligible contribution to the mirror lifetime.

A typical EUV mirror can consist, e.g., of a layer pair of molybdenum and silicon with the molybdenum layer 2.76 nm thick, such that 200 pairs for sacrificial erosion gives, e.g., 552 nm of molybdenum erosion before end-of-life for this mirror. For a 1-year useful life, the molybdenum sputter rate must be below 552 nm/year, i.e., 1.75×10^{-5} nm/sec.

The lithium sputter rate in terms of atoms per cm^2 per second (equal to the lithium influx rate derived above) converts to nm/sec from the thickness of a monolayer of lithium, given the atomic number density of lithium per its mass density and atomic weight, with appropriate unit conversions, as follows:

$$\text{AtomicNumberDensity} = \frac{\rho_{\text{lithium}} (\text{g}/\text{cm}^3)}{M_{\text{lithium atom}} (\text{amu}) \cdot \frac{1.6605 \times 10^{-24} \text{ g}}{1 \text{ amu}}} \{12\};$$

where $\rho_{\text{lithium}} = 0.535 \text{ g}/\text{cm}^3$ and $M_{\text{lithium atom}} = 6.941 \text{ amu}$. The resulting atomic number density for lithium is $4.64 \times 10^{22} \text{ atoms}/\text{cm}^3$. If this number of lithium atoms were arranged in a cube with dimensions 1 cm on each side, then the number of atoms along an edge per cm would be the cube root of the atomic number density, $3.58 \times 10^7 \text{ atoms}/\text{cm}$. The resulting monolayer thickness is $2.78 \times 10^{-8} \text{ cm}$ or 0.278 nm. The number of atoms per cm^2 in a monolayer then is the square of the number of atoms along an edge per cm: $1.28 \times 10^{15} \text{ atoms}/\text{cm}^2$.

The number atoms of, e.g., lithium, removed by sputtering per second must match the influx rate given in Equation 11. Thus, the number of monolayers removed per second is equal to the influx rate divided by the number of atoms per cm^2 in a monolayer. Thickness removal rate is the monolayer removal rate times the thickness of a monolayer, i.e.,

$$\text{ThicknessRemovalRate} = \text{MonolayerThickness}(\text{nm}) \cdot \frac{\text{InfluxRate}(\text{atoms}/\text{cm}^2\text{s})}{\text{NumberOfAtomsinaMonolayer}(\text{atoms}/\text{cm}^2)} \{13\}.$$

Using the values for lithium :

$$\text{Lithium Thickness Removal Rate} = 2.17 \times 10^{-16} \cdot \text{Lithium Influx Rate (atoms / cm}^2\text{ s)} \quad \{14\}$$

with the resulting units of nm/sec. The lithium influx rate shown in FIG. 11 converts to a required lithium thickness sputter rate, shown in FIG. 12, for the same 1-4 droplet sizes, repetition rate and duty cycle

. This result further highlights the need for a small droplet size and a large mirror radius. Otherwise, the required sputter rate can become impractical.

The required thickness sputter rate for lithium, can be compared to the maximum allowed thickness sputter rate for molybdenum, e.g., for a 1 year collector lifetime. The data in FIG. 12 divided into the maximum allowed molybdenum sputter rate, 1.75×10^{-5} nm/sec is shown in FIG.13 for the same 1-4 droplet sizes, repetition rate and duty cycle

. The question is what is needed to create a molybdenum sputter rate 4 or more orders of magnitude less than the lithium sputter rate. The sputter yield for lithium and molybdenum when attacked by helium ions is discussed, e.g., in W. Eckstein, "Calculated Sputtering, Reflection and Range Values", [citation to publication?] _____, June 24, 2002. This sputter yield data versus ion energy is shown in FIG. 14 along with data for silicon for ion energies of (3) lithium into Mo at $E_{th}=52.7$ eV, (2) helium into Si at $E_{th}=10.1$ eV and (1) helium into Li. As one can see, a properly chosen helium ion energy will result in acceptable lithium sputter yield and essentially no molybdenum sputter yield. A problem can arise, however, from the fact that one cannot control the incident ion energy perfectly. That is, the energy spectrum of incident helium ions is not a delta function. It is the spread of ion energies that must be assessed when determining the differential sputtering between lithium and molybdenum.

There are examples in the literature of RF Induction (RFI) plasmas which create, e.g., an ion energy distribution that is Gaussian shaped with, e.g., a FWHM of 2.5 eV as discussed, e.g., in J. Hopwood, "Ion Bombardment Energy Distributions in a Radio Frequency Induction Plasma," Applied Physics Letters, Vol 62, No. 9 (March 1, 1993), pp 940-942.

The peak of the ion energy distribution can, e.g., be adjusted with proper choice of, e.g., electric field strength and helium pressure. By choosing, e.g., a peak ion energy of 20 eV, the helium ions have high sputter yield for lithium, but have energies safely below that of the molybdenum sputter threshold. In FIG. 15 there is shown a plot of normalized ion energy distribution (1 on the log scale and 2 on the linear scale) centered on 20 eV and FWHM of 2.5 eV along with the sputter yields for (3) lithium, (4) silicon, and (5) molybdenum. One can see that there are very few helium ions with energy above the molybdenum sputter threshold. To determine the sputter rate of molybdenum under these conditions requires calculating the influx of helium ions needed to maintain the mirror surface clean of lithium atoms. A constant sputter yield of 0.2 atoms per ion can be assumed, since the bulk of the distribution of helium ion energies falls within the region of nearly constant lithium sputter yield.

$$\text{HeliumIonInflux(ions/cm}^2\text{s)} = \frac{\text{LithiumInflux(atoms/cm}^2\text{s)}}{\text{SputterYield(atoms/ion)}} \{15\}.$$

Thus, the helium ion density must be 5 times the value of lithium influx density shown for various conditions in FIG. 11.

This helium ion influx expressed in Equation 15 may be considered to be the bare minimum, assuming, e.g., the lithium does not deposit perfectly uniformly. In this event a higher total sputter rate may be required, e.g., to ensure that islands of lithium do not develop. On the other hand, other researchers have shown that the ejection of material from an LPP plasma tends to travel toward the laser source. One can, therefore, e.g., arrange the system such that the laser illuminates the lithium droplet from a direction away from the collector, or through an aperture in the collector mirror that causes much of this debris to not strike the collector mirror. Thus, the total lithium load on the mirror may be reduced from the total theoretical amount striking the mirror.

Knowing the total flux of helium ions and assuming a Gaussian energy distribution with a peak at 20eV and a FWHM of 2.5eV, the integral of a normalized Gaussian distribution is $\sqrt{2\pi\sigma^2}$ where σ^2 gives a variance of the distribution related to the FWHM by:

$$\sigma^2 = \frac{(\text{FWHM})^2}{4 \ln(4)} \quad \{16\}.$$

The integral of a normalized Gaussian then is $\sqrt{\frac{\pi(\text{FWHM})^2}{2 \ln(4)}}$, so that the peak current density of helium ions is given by:

$$\text{PeakHeliumCurrentDensity}(\text{ions}/\text{cm}^2\text{s per eV}) = \frac{\text{HeliumIonInflux}(\text{ions}/\text{cm}^2\text{s})}{\sqrt{\frac{\pi(\text{FWHM})^2}{2 \ln(4)}}} \quad \{17\}.$$

Taking the case of a 25 μm droplet with a mirror radius of 10 cm, the peak helium current density must be 3.38×10^{15} ions/cm²s per eV in order to sputter a total of 1.88×10^{15} lithium atoms/cm²s. This helium current density distribution (1) is plotted in FIG. 16 on a log scale, with (2) silicon sputter density and (3) lithium sputter density, along with empirically determined sputter yield of (4) lithium, (5) silicon, and (6) molybdenum and the product of these functions times the ion current density. A surprisingly beneficial result of this analysis shows that the peak sputter density for molybdenum is 3.5×10^{-205} atoms/cm²s per eV(not shown on graph), an incredibly small value. In fact, even the peak silicon sputter density is more than 3 orders of magnitude smaller than that for lithium.

The integral of these sputter densities over all helium ion energies gives the total sputter rate. These integrals are shown respectively as dashed curves (1) for lithium and (2) for silicon, in FIG. 17. The integrated lithium sputter density is 1.88×10^{15} atoms/cm²s, matching the lithium influx rate. The integrated silicon sputter density is 9.17×10^{10} atoms/cm²s. The integrated molybdenum sputter density is 1.16×10^{-205} atoms/cm²s. Therefore, differential sputter rates between molybdenum and lithium are so low that, e.g., less layers of the collector mirror need be employed, e.g., many less than a previously anticipated 300 base pair mirror concept. A single molybdenum layer will last more than a year under these conditions and the assumptions of this sputter yield model. This performance could be even more improved using a debris shield between the ignition spot and the collector main mirror or main an secondary mirrors, but the debris shield, as seen

from these results, may also be totally eliminated, at least for a lithium target. This type of stimulated plasma induced ionized sputtering of debris from the EUV optics, especially for a lithium target, as seen from the above, could even allow for use of other target types, e.g., a moving tape or other type of moving solid target system.

- 5 Helium ion sputtering can be arranged such that it removes the lithium atoms from the collector mirror at a sufficient rate while sputtering molybdenum at a low enough rate for far greater than 1 year lifetime.

Sputtering of molybdenum by, e.g., lithium ions must also be considered in the embodiment of the present invention being discussed, since, e.g., there will be
10 lithium ions formed a debris from the ignition plasma which do not reach the optic surface, but which will be available to the sputtering plasma and will be accelerated toward the mirror surface with a similar energy distribution as the helium ions. The literature also provides data on sputter yield of lithium and molybdenum with lithium ions. This data is shown in FIG. 18 in curve 1 for lithium at $E_{th}=36.3$ eV,
15 along with the same normalized lithium ion energy distribution as was used for the helium ions. To calculate the molybdenum sputter density from lithium the total lithium ion influx must be known. Unlike this calculation for helium (Equation 15) it is not clear what the total lithium influx will be, however, a conservative choice would be the total lithium atomic influx generated by the LPP ignition plasma. Using
20 Equation 17 and the assumptions of a 25 μm droplet and a 10 cm mirror radius, 1.88×10^{15} lithium ions/ cm^2s would be incident on the mirror, and the peak lithium ion current density is 7.06×10^{15} lithium ions/ cm^2s per eV, with the assumption of a 2.5eV FWHM spread in incident ion energy, which, when multiplied by the sputter yield for molybdenum and integrated over all ion energies, gives a total
25 molybdenum sputter density of 2.54×10^{-48} atoms/ cm^2s . This is much higher than that for helium ions, but still much, much lower than the rate required for one year of useful life.

The molybdenum sputter density with lithium ions can be converted to thickness loss rate by using Equations 12 and 13. For molybdenum:

30 $\rho_{\text{moly}} = 10.2 \text{ g/cm}^3$
 $M_{\text{moly atom}} = 95.94 \text{ amu} = 1.59 \times 10^{-22} \text{ g}$

Moly Atomic Number Density = 6.40×10^{22} atoms/cm³

Moly Monolayer Thickness = 2.50×10^{-8} cm = 0.250 nm

Moly Monolayer Atomic Density = 1.59×10^{15} atoms/cm²

Thus, the sputter thickness rate for molybdenum, when attacked by lithium
5 atoms, is 3.99×10^{-64} nm/sec or 1.25×10^{-56} nm/year. This also leads to the conclusion
that the above noted beneficial results of the sputtering plasma ionized cleaning of
the EUV optics by, e.g., helium ion sputtering are still realizable even with, e.g.,
lithium sputtering of molybdenum.

An additional beneficial result is the reconsideration of the previously
10 proposed use of, e.g., a ruthenium capping layer on, e.g., the multi-layer mirror. A
ruthenium capping layer has been proposed to prevent EUV-assisted oxidation of the
first silicon layer in the Mo/Si stack. Multi-layer mirrors are usually terminated with
a silicon layer rather than a molybdenum layer because the molybdenum layer
would quickly oxidize once exposed to room air. Applicants, before the above
15 analysis regarding sputtering plasma cleaning of the EUV optics had considered,
e.g., a multi-layer mirror terminated with silicon, with the expectation that the first
layer of silicon would be eroded to expose the first layer of molybdenum or a
ruthenium capping layer to avoid oxidation of a first layer of molybdenum if that
approach was taken. The super-slow erosion rate of molybdenum, and a similar
20 expected low erosion rate for ruthenium allows for use of a ruthenium capping layer
expected to last for the useful life of the mirror. This results in no loss of the first
layer of silicon, and no need to worry about what havoc the sputtered silicon atoms
might cause, and no oxidation problems with an exposed molybdenum layer. The
sputter yield of ruthenium with lithium and helium, although expected to be similar
25 to that of molybdenum, since ruthenium has a higher atomic mass than
molybdenum, remains to be determined.

The minimum RF power needed to create the desired sputtering plasma at or
near the optic surface can be calculated by assuming, e.g., that every helium ion that
is created strikes the collector mirror, which will underestimate the required RF
30 power, but should give an order of magnitude estimate. Each helium ion that strikes
the collector mirror requires 24.5 eV to ionize, and according to the above example

of an embodiment of the present invention has to have an average kinetic energy of 20eV when it reaches the collector mirror. These two energy values times the required influx of helium ions, 9.40×10^{15} ions/cm²s from Equation 15, gives the plasma power. Converting energy units from eV to J gives a minimum plasma power density of 66.9 mW/cm². Multiplying by the half the surface area of the 10cm radius mirror, 628 cm², gives 42W of minimum total plasma power. Assuming conservatively that only 1% of the plasma power is effectively used, the required plasma power calculated is 4.2kW, which is acceptable, especially considering the very large area over which this power can be dissipated. This estimate of plasma power compares to the previous assumptions of 400mJ per pulse at 6kHz LPP laser power, 2.4kW of laser power and assuming the collector mirror subtends π steradians, it will be exposed to half of this laser power, i.e., 1.2kW. The thermal load from the LPP is similar to the thermal load of the plasma cleaning. The sum of the two powers is 5.4kW, resulting in a power density on the mirror of 8.6W/cm². Applicants believe that a collector mirror exposed to a 10W/cm² or less power density is easily cooled, e.g., with water channels along the back of the mirror, or between the grounded shield and the mirror.

If the plasma power effectivity is more like 10%, then the total power density onto the mirror is only 2.6W/cm², making it possible to radiatively cool the mirror, according to Stefan's law of radiation, which states that the power radiated per square meter from a black body at temperature T is given by:

$$P = 5.67 \times 10^{-12} \cdot T^4 \quad \{18\};$$

where temperature is in Kelvin and the resulting power density is in W/cm², which is plotted in FIG. 19. A temperature in excess of 500°C would be required to radiate all of this incident power, so active cooling of the collection mirror appears to be required in order to prevent damage to the multi-layer stack.

Turning now to FIG. 20 there is shown an schematically an apparatus and method according to an embodiment of the present invention for reclaiming damaged EUV optics, e.g., those that have lost reflectivity, e.g., due to deposition of material on the reflective surface, e.g., carbon and/or carbon based molecules, which may come from, e.g., contamination entering the EUV plasma chamber or from

sputtering or photonic removal from layers of the multi-layer reflective stack coated on reflecting surfaces on the EUV apparatus. As can be seen in FIG. 20 a photo-chemical cleaning apparatus 400 may include a chamber, within which may be mounted; e.g., a collector holding jig 402 that is adapted to hold a collector for
5 cleaning. Also included may be, e.g., a source of photonic energy, e.g., a DUV light source 410, with the collector holding jig 402 and the light source 410 arranged so that the light from the light source 410 simulates light coming from a point source at the focus of the collector, e.g., the ignition site 28 discussed above, such that the collector 404 is irradiated as if by light from a target ignition site.

10 According to an embodiment of the present invention, e.g., the chamber 401 may first be purged by the use of nitrogen provided to the chamber through N₂ valve and then evacuated from the chamber 401 using gas exit valve, followed by the introduction of a fluorine containing gas, e.g., molecular F₂ or NF₃. The collector 404 may then be subjected to irradiation by the light source, e.g., DUV light at a
15 range of λ between, e.g., 160-300nm, e.g., from a KrF excimer laser at 193nm, e.g., in a MOPA configuration for high power at about 40 W, with a pulse repetition rate at about 4kHz. This can serve, e.g., to stimulate the production of, e.g., fluorine based carbon materials, e.g., CF₄, e.g., in a gas phase, which can then be evacuated from the chamber 401 through the gas exit valve 420 under a second nitrogen purge.

20 An alternative of a KrF DUV light source could be, e.g., a commercially available DUV lamp, e.g., a KrCl DUV lamp.

Applicants expect that thicknesses of about 3.5nm carbon atom deposition on an EUV optic, e.g., a collector reflective surface can reduce reflectivity by about 5% and a 10nm deposition by about 14%. Such levels of thickness of deposit are
25 expected to be removed from, e.g., the collector optics reflective surfaces under treatment in fluorine with selected concentrations and the above referenced level of DUV light for a selected time. The process could also employ replenishing the fluorine supply with a gas flow control valve (not shown) to maintain, e.g., a desired concentration of fluorine during the cleaning process.

30 Applicants herein also propose according to an aspect of an embodiment of the present invention that other types of barrier materials may be used in the multi-

layer reflecting mirror stacks to help improve the thermal stability and reflectivity of, e.g., Mo/Si reflective stacks, e.g., optimized for 13.5 nm EUV light reflectivity. To promote smoothness of very thin, e.g., 1nm barrier layers, that are compatible with, e.g., Mo/Si and perhaps also MoSi₂, retaining the appropriate levels of
5 transparency to, e.g., 13.5 nm light, applicants propose the use of inter-diffusion barrier layers comprising carbides selected from the group comprising ZrC, NbC, SiC, borides, e.g., selected from the group ZrB₂, NbB₂, disilicides selected from the group comprising ZrSi₂, NbSi₂ and nitrides BN, ZrN, NbN and Si₃N₄. Other such layers could include yttrium, scandium, strontium compounds and/or these metals in
10 pure form. Among the above, the carbides and borides mentioned are preferred due to the ability to create smoother diffusion barrier layers with such materials.

According to aspects of an embodiment of the present invention applicants contemplate multi-layer stacks, including e.g., MoSi₂/Si, Mo₂C/Si, Mo/C/Si/C and Mo/X/Si/X, where the first two are MLMs where MoSi₂ or Mo₂C is used in place of
15 the Mo normally used in normal Mo/Si mirror coatings, with no inter-diffusion barriers. The other two are with the so-called inter-diffusion barriers, where C refers to carbon and X refers to a suitable material, including further compounds, e.g., the above noted borides, disilicides, and nitrides as the X materials. Nitrides are currently preferred embodiments according to applicants for inter-diffusion barrier
20 layers in the applications according to embodiments of the present invention. Mo₂Si/Si is described in the paper Y. Ishii et al. "Heat resistance of Mo/Si, MoSi₂/Si, and Mo₅Si₃/Si multilayer soft x-ray mirrors", J. Appl. Phys. 78, (1995) p. 5227.

Helium has high transparency to EUV, which makes it a good choice for a
25 buffer gas for which a transmission of 90% is representative. Based on the partial pressures required for efficient sputtering, a few mTorr, helium buffer gas transmission would be nearly 100%. A possible collector multi-layer surface could comprise, e.g., 300 coating pairs instead of the normal 90 pairs. The extra pairs would not improve the reflectivity over a 90 pair mirror, but instead these extra
30 layers can, if required, get used once the top layers are eroded away. With a 300 pair mirror the sputter rate differential between lithium and the mirror need not be so

high that a single mirror layer lasts for, e.g., months at a time. Instead three could be, e.g., an extra 210 layer pairs worth of mirror erosion that can be sustained.

Lithium chemical compounds that might be generated in the LPP vessel, e.g., LiH, LiOH, Li₂CO₃, etc., can have melting points in excess of 600°C and thus may not be evaporated from the mirror. These could even form in certain cases, e.g., a crust over the lithium which deposits on the mirror surfaces. These could, however, very effectively be sputtered by the sputtering ion plasma, e.g., containing the ionized He atoms, or may be sputtered by lithium itself in the form of high speed lithium ions and atoms ejected from the plasma that impinge on the reflecting surface.

The sputter rate required to stay ahead of the lithium deposition could be much higher in an EUV light source than the literature indicates is typically practiced, e.g., in modern deposition and etch machines, which is at least part of a reason for, e.g., a combined approach to keeping the, e.g., lithium off of the mirror surfaces. According to an aspect of an embodiment of the present invention applicants contemplate using evaporation to remove the bulk of the lithium while employing a very light sputter rate to remove the inevitable lithium and carbon compounds deposited on the mirror surface. However, even a very light sputtering plasma impinging on at least the main and secondary reflecting surfaces could have the same beneficial carbon and other lithium compound removal properties. Employing this idea beyond the intermediate focus, e.g., in the illuminator reflective surfaces and also the projection reflecting surface may also prove beneficial to remove debris that happens to reach the lithography tool reflecting surfaces. In the lithography tool itself, due to, e.g., smaller deposition rates the thermal load and sputtering rate may be sufficiently low for this to be effective.

Sputtered lithium and lithium compounds along with lithium ejected from the plasma that does not collect on the reflecting surface may be trapped in cold fingers [not shown] contained in the EUV light source vessel, e.g., in the form of cooled, e.g., water cooled fins or plates extending from the inside walls of the vessel, and out of the optical path from the collector to the intermediate focus.

In the case of, e.g., tin as the source element it may be possible to use, e.g., a hydrate of the metal, e.g., SnH_4 , which is a vapor at room temperature, along with a hydrogen plasma for cleaning the collector in a tin-based LPP source. Hydrogen has high 13.5nm transmission and the resulting SnH_4 could be pumped away rather than trapped on cold fingers like the lithium.

Applicants have examined, e.g., the stopping power of helium and argon against both tin and lithium ions. The results are shown in FIG. 21A and B. The two graphs have the same data, just different scales. Lines 500, 502 and 503, for tin at different measured distances from a source plasma, respectively 96.5 cm, 61 cm and 32.5 cm with solid being helium buffer and dashed being argon buffer. The lines 506 are for lithium. If pressure*distance product scaling were applied, these three sets of for tin data would fall approximately on top of each other.

Applicants have also determined that Argon has at least 10 times higher stopping power than helium for a given gas pressure. Also, lithium can be stopped with less buffer gas than tin. And, scaled to the true working distance of an LPP collector (~10cm), the required buffer pressure, even with argon, will need to be in the range of about 10mT for tin. Since xenon and tin have nearly the same atomic mass, applicants expect that the required buffer pressure for a xenon LPP would also be in the range of 10mT. Such a high buffer gas pressure can present EUV self-absorption problems for xenon and tin. But not for lithium, both because of the lower buffer pressure requirement and also the lower EUV absorption of lithium.

In continuing to examine the stopping power of a buffer against, e.g., the fast ions produced by the LPP, using, e.g., a Faraday cup to collect and measure the ions at a known distance through a known aperture size at different increasing buffer gas pressure this Faraday cup signal decreased, giving a measure of the ion stopping power. The results for tin and lithium are shown below in FIG.'s 22A-E. FIG.'s 22A and 22B show the raw Faraday cup signal vs. time for, respectively tin and lithium. In FIG.'s 22C and D these signals are plotted vs. ion energy using time-of-flight respectively for tin and lithium. In FIG. 22E the area under these curves is plotted vs. a pressure*distance product of the buffer gas, with the lower plot line (1) being tin and the upper (2) being lithium.

A surprising result of this analysis by applicants was that the last graph shows the Faraday cup signal vs. buffer gas $P \cdot D$ product for both tin and lithium being about the same for both elements. Applicants believe this is explainable in that the analysis was not really measuring the loss of ions captured by the Faraday cup, but instead was measuring the neutralization of the ions by the buffer gas, so-called electron capture by the ions. If an ion is neutralized, it will not register in the Faraday cup. This can be explained, e.g., because tin might have a larger electron capture cross-section than lithium, especially considering that the tin ion is highly charged, 7-11 times ionized and the lithium can, at most, be 3 times ionized. The stopping power result shown in FIG. 22E can be considered an overestimate of the buffer gas stopping power in that it can be no better than the value predicted by these curves.

Taking the observed values of the stopping power as the upper limit one can calculate the necessary pressure of argon buffer gas to extend the collector mirror lifetime to 100B pulses. Starting with the result from Engineering Test Stand (ETS) built by the EUV LLC, which reported that one multi-layer mirror pair is eroded for every 15M pulses with a xenon LPP and a collector distance of 12cm, and assuming that the reflectivity of a multi-layer mirror is not significantly degraded until 10 layer pairs are removed, the ETS collector mirror had a lifetime of 150M pulses compared to a requirement of 100B pulses. This leads to the conclusion that a reduction of 666X in erosion rate is necessary. On the plot in FIG. 22E a $P \cdot D$ product of approximately 500mT*cm would be required to achieve this level of reduction. A working distance of, e.g., 12cm gives, e.g., a need for an argon pressure of 42mT. This also can result in the conclusion that lithium is the better target over, e.g., xenon, since for xenon LPP, a buffer pressure of 42mT is not very satisfactory due to the strong EUV absorption of the xenon caught up in the argon buffer. For lithium, however, this amount of buffer pressure is no problem for lithium absorption. Tin may also be satisfactory, depending on, e.g., the vapor pressure and evolution rate of SnH_4 from the surface of the collector mirror. A relatively large buffer gas pressure seems, therefore, to be a requirement, which

leads to the conclusion that xenon is not a good target, tin may be, but lithium appears to be the best.

Applicants have also determined that even if the effectiveness of heating the collector reflective surfaces is impacted by the fact that the material being evaporated needs to have, e.g., a certain thickness, e.g., 50 Å, e.g., about 10 monolayers, before published values of vapor pressure are realized, i.e., the material, e.g., lithium may be harder to evaporate directly off of the surface of the mirror, nevertheless, the transmittance of such a thickness of lithium on the mirror surfaces is about 95%, and about 90%, double-pass, so that such a layer on the mirror would not significantly detract from the overall CE, e.g., at 13.5nm. In addition, such a layer of “evaporationless” lithium, may actually be beneficial in that it may be able to protect the collector mirror from the onslaught of high-speed lithium atoms and ions. This lithium layer will be sputtered instead of the molybdenum layer of the multi-layer mirror. Xenon, because it is a gas, will not form such a protective layer and a tin layer, because of its very high EUV absorption, would only be 52% transmitting.

Given that the sputter yield of lithium against molybdenum is much less than the sputter yield of xenon against molybdenum, e.g., for ion energies around 1keV (the sputter rates tend to saturate above this energy level):

	Incident Ion		
	<u>Target Material</u>	<u>Lithium</u>	<u>Xenon</u>
	Lithium	0.21	??
	Molybdenum	0.081	1.45

xenon will sputter molybdenum at 18 times higher rate than lithium. This difference alone would give a 2.7B pulse collector lifetime without changing anything else. The “evaporationless” steady state thin layer of lithium may provide the remaining 37X reduction in sputter rate. Even if it does not, the EUV LLC concept of producing a mirror with ~100 extra sacrificial layer pairs could add a 10X increase

in lifetime, e.g., to 27B pulses, which combined with the lower erosion from lithium could give a collector lifetime of 100B pulses.

Applicants have also examined the effectiveness of electrostatic protection of the collector mirror. The concept has been proposed in the literature, i.e., to
5 generate an electric field between the source LPP and the collector mirror such that the energetic ions must climb up a potential well as they travel toward the mirror. This potential well can be made deep enough that the ions lose all of their kinetic energy before reaching the mirror. In fact, they are turned around and sent packing back down the potential well, never reaching the mirror. Applicants have
10 discovered, however, that attempting to do this by running an electrical connection through the vessel to the collector mirror was ineffective due to the target bias dropping to near zero upon pulsing the laser, which was determined to be the result of a high peak current required to maintain the bias voltage and the large lead wire required, thus dropping all of the voltage along the inductance of the wire. To
15 correct this problem applicants then installed capacitors inside the vacuum vessel and constructed low inductance buss work between ground and the target plate. Inductance was measured by placing a copper sheet up and around the target and attached to ground. By charging the capacitors to a low voltage and discharging them by pressing the copper sheet against the target applicants measured the ringing
20 voltage waveshape and inferred the inductance. The result was 104nH with a 697ns half-period discharge waveshape. This discharge period is much longer than the laser pulse and subsequent EUV emission initially causing concern whether the bias could be maintained during the critical period when the ions are created and leave the plasma region (~20ns). Applicants determined, however, that such short time
25 scales were unimportant. What is important is, e.g., to maintain or reestablish the target bias, e.g., in a time scale that is, e.g., short compared with the travel time of the ions from the target to the mirror. With the present geometries the ion travel time is about 2.5 μ s, so a circuit half-period of 0.7 μ s should be sufficient.

In testing this arrangement applicants were surprised to find that the full
30 0.47 μ F capacitance was drained of its -1000V potential in a time scale almost exactly the same as during the inductance measurement using the copper strap.

Applicants determined that the laser pulse initiates a discharge between the target plate and the vessel wall. This discharge completes the circuit between the capacitor's high voltage terminal and ground, thus draining the capacitors as if a copper strap had been placed across them. Evidently, the events that unfold during, and immediately after, the laser pulse, a plasma being created at the target point and this plasma radiating a large amount of hard UV and EUV radiation throughout the vessel. The energy of most of these photons is above the work function of the metals inside the vessel and thus photoelectrons are created at all the metal surfaces. These photons are also energetic enough to ionize any gas atoms that exist in the vessel. In this case argon was used as the buffer gas and it is easily ionized by the hard UV and EUV radiation produced by the LPP. And finally, electrons and ions are created in the LPP and stream outward into the volume of the vessel. Except for those ions that are attracted to the biased target plate. They strike the plate and create secondary electrons. Essentially, the creation of a discharge between two metal plates held at a potential between each other occurs as if the arrangement were a laser-triggered discharge switch.

There still is some possibility of making an effective electrostatic repulsion, but it becomes a bit more complicate and isn't really electrostatic. The idea is to pulse the bias such that it is present only after the initial events of the laser pulse. In only a few 100's of ns most of the electrons will have collided with the vessel wall, and of course the radiation will be gone. At this time it might be possible to apply a bias and repel, or attract, the ions away from the collector mirror.

Those skilled in the art will appreciate that the above references preferred embodiments of the present invention and aspects thereof are not meant to be exclusive and other modifications and additions to the above referenced embodiments may be made without departing from the spirit and scope of the inventions disclosed in the present application. The appended claims, therefore, should not be considered to be limited to the above disclosed embodiments and aspects but should include with the scope and spirit of the claims the recited elements and equivalents of the recited elements. By way of example, other target material and multi-layer reflective coating metals may have similar relationships as

discussed above to allow for the continuous cleaning by, e.g., sputtering, e.g., of ions, e.g., induced by the creation of a sputtering plasma in the vicinity of the optic reflecting surface(s), which ions may also be, e.g., other than helium, e.g., H, N or O. Also, e.g., the heating mechanism for the reflecting surface could be a heat lamp
5 directed at the reflective surface. Other such changes and additions may be appreciated by those skilled in the art.